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IR-Spectroscopic Investigations of Cycloalkanes*

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Summary

A homologous series of cycloalkanes from CH_2)₁₄ up to $(CH₂)$ 96 as model systems for tight folding in CH₂-chain molecules was investigated by infrared spectroscopy. In comparison with the corresponding spectra of the nalkanes we found two additional new absorption bands at 700 cm ' and 1442 cm⁻' which can be assigned to a characteristic fold vibration. The already known fold band in the wagging region at about 1344 cm⁻¹ could also be found in all cycloalkane spectra. A numerical decomposition of the bands enables the calibration of the fold concentration. The Oavydov splitting of the (CH2)-rocking and ICH2)-bending vibration in the melt cry stallized form of the molecules (CH2)72 and (CH2) $q₆$ show the occurence of a orthorhombic subcell of the larger rings whereas the smaller rings and the larger ones in the solution crystallized form only crystallize in a monoclinic form.

Introduction

In order to investigate the influence of tight chain folding in polymethylene chains on the spectroscopic behaviour it is useful to have an appropriate model system. Cyeloalkanes represent such a system because this molecules crystallize as double stems linked by two folds (NEWMAN and KAY, 1976; TRZEBIATOWSKI and STROBL, 1980). From semiempirical atomistio calculations with a proved potential set one finds three different fold conformations for PE with rather low defect energies (fig.l). Because of the low energy and the good packing of the b-axis fold (fig.la) it is a reasonable assumption that this type of conformation (..ttggtggtt..) will also form the folds in crystals

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Fig. 1 a,b,c: Possible fold conformations in PE a} b-axis fold, defect energy 16.3 kJ/mole, b) diagonal fold, defect energy 20.4 kJ/mole, c) fold in a CH₂-chain molecule existing as $4 \times 1/1$ helix, defect energy 13.4 kJ/mole, (GROSSMANN, 1977)

which are grown from solution. In the melt crystallized material of the larger ring molecules the diagonal fold (fig.lb) may be favoured because of the orthorhombic position of the straight chain sequences between the $folds. Fig. 2 shows a model of the $(CH_2)_{30}$ -molecule$ formed by parallel trans segments linked by two tight folds.

Fig. 2: Possible conformation for a $(CH₂)₃₀$ -molecule in the solid state.

Experimental

The IR-spectra were measured with a grating spectrometer model 283 from Perkin Elmer. Because of the limited quantity of specimen material some of the experiments were done with KBr moulding technique. For most of the quantitative analysis of the absorption bands we used cells with KBr windows. The temperature dependent measurements were done with a self-made cell for liquids which is temperature controlled by an

Fig. 3: Spectra of the rocking, wagging and bending vibrations for some cycloalkanes. ((CH $_{\rm 2}$) $_{\rm 72}$ and (CH $_{\rm 2}$)g $_{\rm 6}$ monoclinic modification) (ARNOLD, 1979, see also TRZEBIATOWSKI et al., 1980).

extern liquid thermostat. The cycloalkanes were synthesized by SCHILL and coworkers with a special method reported some years ago (SCHILL et al., 1976).

Results and discussion

IR-spectra: Fig. 3 shows the IR-spectra of the eycloalkanes in the region of the rocking vibrations and the spectra in the wagging and bending region. A comparison of the band positions with the vibrational analysis by KRIMM et al. {KRIMM and JAKES, Ig71) for (0H2)34 1 support the assignment 1368 cm¯' (gtg), 1352 cm¯' (gg and contributions of other g-conformations), 1344 cm⁻¹ (ggtgg). These bands could be found in all cycloalkanes investigated. The intensity of the 1344 band with respect to the rocking mode decreases with increasing molecular weight. If we compare the IR-spectra of the n-alkanes with those of the ring molecules (fig.3, fig. 4) we can establish that there appear two additional bands. Une of them is proposed to be a rocking CH₂-mode, located at 700 cm⁻' exactly. The other one, located at 1442 cm^{-1} , is assumed to be a CH₂-bending vibration. The position of these two additional bands is unaltered for all cycloalkanes investigated whereas the intensities of the bands decrease with increasing number of CH₂-groups per molecule. This is why in the spectra of the larger molecules the bands are no longer peaks but are reduced to shoulders of the rocking and the bending vibration, respectively.

If we look at the spectra of the $(CH_2)_{72}$ and $(CH_2)_{96}$ we find differences between the melt crystallized (fig.4) and the solution grown form. In the melt crystallized material the spectra of these molecules show the characteristic Davydov-splitting of the rocking and bending vibration. This indicates a orthorhombic $crystal line$ subcell. The asymmetry of the band splitting - more pronounced for $(CH_2)_{72}$ than in the $(CH_2)_{96}$ spectrum - shows that the melt crystallized material contains already monoclinic or triclinic parts besides the orthorhombic crystallites. The existence of the orthorhombic components can be also seen in the corresponding raman-spectra of $(CH_2)_{72}$ and $(CH_2)_{96}$ as the occurence of the characteristic band at 1418 cm -I (GROSSMANN and 80LSTLER, 1980; 80LSTLER, 1978).

Band-analysis: In order to obtain a quantitative result the intensities and the integral intensities of the bands (as compared with reference bands at 715 cm⁻¹ (rocking CH_2) and 1472 cm⁻¹ (bending CH_2)) were calculated from the measured spectra. In some cases, especially for the band complex in the wagging region we had to decomposite the single bands numerically. For this purpose we used a computer method (PITHA and JONES, Ig68) which gives the possibility to fit each band to

a Gaussian or a Lorentzian or a mixed shape. Fig.5 gives an example for this method.

Band-intensities: Fig.6 shows the relative absorbance of the bands at 700 cm~' , 1442 cm~' and 1344 cm⁻' of the cycloalkanes at room temperature. The nearly linear dependence on the reciprocal molecular weight as well as the fixed position of these bands for all ring molecules show that these bands must be caused from the tight fold conformation in the molecules. The strong deviation from the linear dependence {fig.6] for the molecules with about n < 28 is certainly caused by **per-**

turbation of the fold structure because of the inner stress of the smaller molecules. The relative integral intensity of the various fold bands shows the same dependence on the molecular weight as the relative absorbance shown in fig.6. The comparison of the IR-spectra in the monoclinic phase with those of the orthorhombic phase (fig.3,4) in [CH2)72 and [CH2)96 shows that the integral intensity of the 1368 cm'' band is greater in the orthorhombic modification. This increase of gig sequences can be understood with a increasing number of diagonal folds (fig.lb) each containing a gig sequence. A detailed consideration of the variation of the defect intensities with temperature will be given elsewhere.

For a rough classification one can establish that the melt transition of the ring molecules (n > 30) causes a strong increase of the 1368 cm~' and 1352 cm⁻' band intensities and a decrease of the intensity at 1344 cm^{-1} . The same rule shall apply for the smaller rings (n < 30) with the restriction that these molecules have a relatively strong absorption at 1368 cm⁻¹ already at room temperature (fig,7), Common to all melt spectra is the appearence of a small band at 1385 cm⁻' which we are net able to assign,

Nevertheless we think we have described all effects of tight chain folding on the IR-spectra of CH₂-chain moleeuIes. Whether they can be found also in the spectrum of polyethylene has to be investigated in further experiments.

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